## **EPILOGUE**

Despite of more than two decades of extensive research on oxygen isotopic fractionation in ozone isotopomers, there are still a number of unresolved problems related to fundamental fractionation mechanism during its formation and dissociation. Moreover, there are several unanswered issues regarding stratospheric ozone and its direct or indirect heavy isotope transfer to other oxygen containing stratospheric trace species.

The present thesis explores the phenomenon of mass independent fractionation in light of the several experimental processes involving ozone and its interaction with other oxygen containing molecules like  $CO_2$  to throw light on some of these issues.

In the course of this work, a number of experiments were devised in a way such that some of the fundamental fractionation mechanisms related to ozone as well as some specific stratospheric issues (e.g. altitudinal variation of enrichment in stratospheric ozone, relative variation of enrichment in <sup>17</sup>O and <sup>18</sup>O of ozone in the upper stratosphere etc.) could be addressed. Some of the important findings of this study are the following:

Dissociation of ozone contributes significantly in the isotopic enrichment of ozone while recycling is allowed to take place during formation through oxygen photolysis at low pressure (< 50 torr). Over and above the dissociation effect, the amount of ozone produced has also a role in enrichment process. The data are explained easily by introducing a parameter called "turn-over time" ( $\tau = O_3$  reservoir amount / rate of  $O_3$  dissociation) which clarifies the role of dissociative enrichment from the perspective of a Rayleigh type of process. An effort is made to explain and predict the altitudinal enrichment variation of stratospheric ozone with the help of this parameter.

It is established that the isotopic fractionations during photo-dissociation of ozone in Hartley (peak around 254 nm) and Chappuis band (peak around 600 nm) are distinctly different. The former shows a mass independent character while the latter is strictly a mass dependent process. Further investigations on Hartley band dissociation decipher the fact that pure UV dissociation is a mass independent process, which proceeds with equal enrichment in <sup>17</sup>O and <sup>18</sup>O in the left-over ozone. An explanation for this dissociation process is presented in the context of Gao-Marcus theory. An interesting phenomenon is observed during ozone dissociation by its interaction with a surface. The origin of the mass independent character of this process is hypothesized by non-statistical breakdown of short-lived complex  $O_3^*$  (formed by adsorption of  $O_3$  in the wall).

During the isotopic exchange between the ozone photolysis product  $O(^{1}D)$  and  $CO_{2}$ , the  $\delta^{17}O$  and  $\delta^{18}O$  of evolved  $CO_{2}$  defines a line of slope 1.7 (with initial  $O_{3}$  and  $CO_{2}$  compositions identical to atmospheric composition) similar to the one observed for stratospheric  $CO_{2}$ . The slope of the line changes with the change of initial  $CO_{2}$  composition and establishes the fact that the isotopic transfer favors <sup>17</sup>O relative to <sup>18</sup>O. It is postulated that a process similar to resonant absorption affects the quenching of  $O(^{1}D)$  such that <sup>17</sup>O containing isotopomers of  $CO_{2}$  is favored during the singlet-triplet transition of the  $CO_{3}^{*}$  complex during its breakdown to O-atom and  $CO_{2}$  molecule.

In summary, this work is an effort to enhance our understanding of the mass independent isotopic fractionation processes.

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# APPENDIX

### Physical Data for Ozone as found in literature:

(Determined by microwave spectroscopy)

Table A1. Physical quantities of ozone molecule.

Quantity	Value
Density (Gas)	2.133 g/L at 273.15 K
Density (Liquid)	1.614 g/cm <sup>-3</sup> at 77.75 K
Melting Point	80.0 K
<b>Boiling Point</b>	161.80 K
Critical Temperature	261.05 K
Critical Pressure	53.8 atm
Critical Volume	$89 \text{ cm}^3/\text{mol}$
Heat of Formation	34.4 kcal/mole (298.15 K)
Heat of Vaporization	4.88 kJ/g
Bond Length	1.2716 Å
Molecular Angle	117.47 °
Point Group	$C_{2v}$



Figure A1. A schematic representation of three types of vibration in ozone molecule.

Isotonomors	Mass	Cal	culated (ci	n <sup>-1</sup> )	Me	easured (cn	n <sup>-1</sup> )
isotopoiners	Iviass	$\nu_1$	$v_2$	<b>v</b> <sub>3</sub>	$\nu_1$	$v_2$	<b>v</b> <sub>3</sub>
$^{16}O^{16}O^{16}O$	48	1103.9	701.3	1043.3	1103.1	700.9	1042.1
$^{17}O^{17}O^{17}O^{17}O$	51	1071.0	608.3	1012.2	1070.9		1012.2
$^{18}O^{18}O^{18}O$	54	1040.8	661.2	983.7	1041.6	661.5	984.8
$^{16}O^{16}O^{17}O$	49	1095.6	692.7	1037.0	1095.7	692.4	1035.7
$^{16}O^{17}O^{16}O$	49	1088.2	697.3	1025.0	1087.8	697.1	1024.4
$^{16}O^{16}O^{18}O$	50	1088.1	685.0	1031.4	1090.4	684.6	1028.1
$^{16}O^{18}O^{16}O$	50	1074.2	693.5	1008.4	1074.3	693.3	1008.4
$^{17}O^{17}O^{16}O$	50	1079.6	688.9	1018.6			
$^{17}O^{16}O^{17}O$	50	1087.2	684.0	1030.7			
$^{17}O^{17}O^{18}O$	52	1063.3	672.5	1006.4			
$^{17}O^{18}O^{17}O$	52	1056.5	676.7	995.4			
$^{18}O^{18}O^{16}O$	52	1057.5	677.7	996.1	1060.7	677.5	993.9
$^{18}O^{16}O^{18}O$	52	1072.3	667.9	1019.4	1072.2	668.1	1019.4
$^{18}O^{18}O^{17}O$	53	1048.7	669.0	989.5			
$^{18}O^{17}O^{18}O$	53	1055.6	664.5	1000.6			
$^{16}O^{17}O^{18}O$	51	1071.9	681.3	1012.9			
<sup>16</sup> O <sup>18</sup> O <sup>17</sup> O	51	1065.4	685.2	1001.9			
<sup>17</sup> O <sup>16</sup> O <sup>18</sup> O	51	1079.8	676.1	1025.1			

**Table A2.** The calculated and measured vibrational frequencies of different ozone isotopomers (taken from Hathorn and Marcus, 2001).

#### Calculated Fractionation of Ozone at Different Wave Lengths Following Miller-Yung Model

The absorption cross-sections of  ${}^{48}O_3$  and  ${}^{50}O_3$  are slightly different which can cause isotopic fractionation due to photo-dissociation. This can be calculated using Miller-Yung model (2000). First, the zero point energy difference ( $\Delta ZPE$ ) between  ${}^{48}O_3$  and  ${}^{50}O_3$  ( ${}^{16}O^{18}O^{16}O + {}^{16}O^{16}O^{18}O$ ) is calculated using the vibrational frequencies of ozone isotopomers, where  $\Delta ZPE$  is given by:

 $(\Delta ZPE)_s = ZPE \ ({}^{16}O{}^{16}O{}^{16}O) - ZPE \ ({}^{16}O{}^{18}O{}^{16}O)$ and,  $(\Delta ZPE)_a = ZPE \ ({}^{16}O{}^{16}O{}^{16}O) - ZPE \ ({}^{16}O{}^{16}O{}^{18}O)$ 

where, 's' and 'a' subscripts denote symmetric and asymmetric case respectively.

Now, ZPE =  $\frac{1}{2}$  ( $v_1 + v_2 + v_3$ ) (where v's are vibrational frequencies of the ozone molecule expressed in cm<sup>-1</sup> unit). Using the vibrational frequencies from Hathorn and Marcus (2001), the calculated  $\Delta$ ZPE values are,

$$(\Delta ZPE)_a = -22.0 \text{ cm}^{-1}$$
  
 $(\Delta ZPE)_s = -36.2 \text{ cm}^{-1}$ 

The <sup>18</sup>O-enrichment in ozone (i.e. the relative difference in absorption rate) can be written as,  $(\sigma_{48} - \sigma_{50})/\sigma_{48} \times 1000$  (‰), where  $\sigma_{48}$  and  $\sigma_{50}$  at a given wave number (v/c) are the absorption cross- sections of <sup>48</sup>O<sub>3</sub> and <sup>50</sup>O<sub>3</sub> respectively.  $\sigma_{48}$  and  $\sigma_{50}$  are related by the equation,

$$\sigma_{50} \left( \nu/c + \Delta ZPE \right) = \sigma_{48} \left( \nu/c \right)$$

#### For ozone dissociation at 253.7 nm:

Using the ozone absorption cross-section data (DeMore et al., 1997) given below, a

function of the fo	orm: $\sigma = a \exp\left(-0.5\left\{\frac{(\nu/c) - 1}{b}\right\}\right)$	$\left \frac{x_o}{x_o}\right ^2$ was fitted.
	$v/c (cm^{-1})$	$\sigma$ (× 10 <sup>-20</sup> cm <sup>2</sup> ) of <sup>48</sup> O <sub>2</sub>

v/c (cm <sup>-</sup> )	$\sigma (\times 10^{-5} \text{ cm}^{-}) \text{ of } {}^{10}\text{O}_3$	
41779.98	797	_
41279.98	900	
40779.94	1000	
40279.88	1080	
39779.86	1130	
39279.86	1150	
38779.90	1120	
38279.85	1060	
37779.78	965	
37279.79	834	
36779.76	692	

The corresponding fit-parameters for <sup>48</sup>O<sub>3</sub> are given below:

Parameters	Values	
а	1147.9544	
b	2684.8433	
X <sub>o</sub>	39386.9557	

Using these parameters,  $\sigma_{48}$  and  $\sigma_{50}$  at 253.7 nm (= 39416.63 cm<sup>-1</sup>) was calculated and we

obtain:	$(\Delta \sigma)_{a} = \sigma_{48} - \sigma_{50} \approx -0.065 \times 10^{-20} \text{ cm}^{2}$
therefore,	$(\delta^{18}O)_a = (\Delta\sigma)_a / \sigma_{48} \times 1000 \approx -0.056$ ‰.
Similarly,	$(\Delta\sigma)_s = \sigma_{48} - \sigma_{50} \approx -0.067 \times 10^{-20} \text{ cm}^2$
and,	$(\delta^{18}O)_{s} = (\Delta\sigma)_{s} / \sigma_{48} \times 1000 \approx -0.058 $ %.

Appendix

finally, 
$$(\delta^{18}O)_{total} = 1/3 \ (\delta^{18}O)_s + 2/3 \ (\delta^{18}O)_a \approx -0.06 \ \%$$

Therefore the left-over ozone will be depleted by 0.06 ‰.

#### For ozone dissociation at 520:

The absorption cross-section data around 520 nm are the following:

$v/c (cm^{-1})$	$\sigma$ (× 10 <sup>-23</sup> cm <sup>2</sup> ) of <sup>48</sup> O <sub>3</sub>
19442.39	162.3
19337.50	173.9
19233.73	182.6
19131.07	191.3
19080.33	205.8
19029.50	217.4

A function of the form:  $\sigma = a \times (v/c)^3 + b \times (v/c)^2 + c \times (v/c) + d$  is fitted to the above data set and obtained the following parameters:

Parameters	Values
a	$-1.32367975047 \times 10^{-6}$
b	$7.65896904133 \times 10^{-2}$
с	- $1.47725859141 \times 10^3$
d	$9.49838314502 \times 10^{6}$

Using these parameters,  $\sigma_{48}$  and  $\sigma_{50}$  at 520 nm (= 19230.77 cm<sup>-1</sup>) was calculated and we

obtain:	$(\Delta \sigma)_a = \sigma_{48} - \sigma_{50} \approx -1.90 \times 10^{-23} \text{ cm}^2$
therefore,	$(\delta^{18}O)_a = (\Delta\sigma)_a / \sigma_{48} \times 1000 \approx -1.90 / 181.9 = -10.45 $ %.
Similarly,	$(\Delta \sigma)_{\rm s} = \sigma_{48} - \sigma_{50} \approx -3.28 \times 10^{-23}  {\rm cm}^2$
and,	$(\delta^{18}O)_s = (\Delta\sigma)_s / \sigma_{48} \times 1000 \approx -3.28 / 181.9 \times 1000 = -18.05 $ %.
finally,	$(\delta^{18}\text{O})_{\text{total}} = 1/3 \ (\delta^{18}\text{O})_{\text{s}} + 2/3 \ (\delta^{18}\text{O})_{\text{a}} \approx -13.0 \ \%$

Therefore, during photo-dissociation of ozone at 520 nm, the left-over ozone will be depleted by 13.0 ‰.

#### For ozone dissociation at 630 nm:

$v/c (cm^{-1})$	$\sigma$ (× 10 <sup>-23</sup> cm <sup>2</sup> ) of <sup>48</sup> O <sub>3</sub>
16048.53	385.5
15976.99	373.9
15906.09	359.4
15836.07	342
15766.4	330.4
15731.68	315.9

The absorption cross-section data around 630 nm are the following:

A function of the form:

$$\sigma = a \times (v/c)^{5} + b \times (v/c)^{4} + c \times (v/c)^{3} + d \times (v/c)^{2} + e \times (v/c) + f$$

is fitted to the above data set and obtained the following parameters:

Parameters	Values
a	$3.24118687744 \times 10^{-10}$
b	- $2.57631099424 \times 10^{-5}$
с	$8.19118920805 \times 10^{-1}$
d	- $1.30214725893 \times 10^4$
e	$1.03499373478 \times 10^8$
f	- $3.29055995954 \times 10^{11}$

Using these parameters,  $\sigma_{48}$  and  $\sigma_{50}$  at 630 nm (= 15873.02 cm<sup>-1</sup>) was calculated and we obtain: ( $\Delta \sigma$ )<sub>a</sub> =  $\sigma_{48} - \sigma_{50} \approx 5.36 \times 10^{-23} \text{ cm}^2$ therefore, ( $\delta^{18}\text{O}$ )<sub>a</sub> = ( $\Delta \sigma$ )<sub>a</sub> / $\sigma_{48} \times 1000 \approx 5.36$  / 359.2 = 14.9 ‰. Similarly, ( $\Delta \sigma$ )<sub>s</sub> =  $\sigma_{48} - \sigma_{50} \approx 8.33 \times 10^{-23} \text{ cm}^2$ and, ( $\delta^{18}\text{O}$ )<sub>s</sub> = ( $\Delta \sigma$ )<sub>s</sub> / $\sigma_{48} \times 1000 \approx 8.33$  / 359.2 × 1000 = 23.18 ‰.

finally, 
$$(\delta^{18}O)_{\text{total}} = 1/3 \ (\delta^{18}O)_{\text{s}} + 2/3 \ (\delta^{18}O)_{\text{a}} \approx 17.7 \ \%$$

Therefore, during photo-dissociation of ozone at 630 nm, the left-over ozone will be enriched by 17.7 ‰.

# **PAPER PUBLISHED/ SUBMITTED**

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## CONFERENCE PROCEEDINGS

- 1. Chakraborty, S., and S.K. Bhattacharya, Oxygen Isotopic Anomaly in Ozone Dissociation on Glass Surface, *Proceedings of Ist International Symposium on Isotopomers (ISI 2001)*, Yakohama, Japan, ISI2001-15.pdf.
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#### **ABSTRACTS**

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